Control of Spin Systems



The Nuclear Spin Sensor

- Many Atomic Nuclei have intrinsic angular momentum called spin.
- The spin gives the nucleus a magnetic moment (like a small bar magnet).
- Magnetic moments precess in a magnetic field at a precession frequency that depends on the magnetic field strength.
- The spins are therefore beautiful, very localized (angstrom resolution) probes of local magnetic fields.
- The chemical environment of a nucleus in a molecule effects the local magnetic field on the nucleus.



• Probing spins with radio frequency magnetic fields and observing them gives information about chemical environment of the nuclei in a non-invasive way. Therefore field of nuclear magnetic resonance is the single most important analytical tool in science.



- The magnetic moment of a single nuclear spin is too weak to detect.
- The spins are generally detected in the Bulk by making them precess coherently.
- The precession of nuclear spins is detected or observed by a Magnetic resonance technique.









Chemistry Pharmaceuticals Life Science Imaging Food Material Research Process Control









Marx et al. (2000)

NMR spectrometer

superconducting magnet incl. probe (with rf coil)



One dimensional spectrum





 $s(t) = \eta \exp(-R_{I}t_{2}) \exp(-R_{s}t_{1}) \cos(2\pi v_{s}t_{1}) \cos(2\pi v_{I}t_{2})$

Example: ¹⁵N-HSQC of p63



¹⁵N labeling:

- all N atoms replaced by 15 N (ca. 95 % 15 N),
- characteristic fingerprint spectrum
- p63: 233 a.a. / 27 kDa
- measured at 750 MHz / 303 K $\,$

primary structure (amino acid sequence)



Relaxation Optimized Coherent Spectroscopy Singular Optimal Control Problems

Transfer of Polarization



Random collisions with solvent molecules causes stochastic tumbling of the protein molecules



$$\frac{1}{\pi T_2} = k \approx J(0)$$

 $T_1 \gg T_2$

Spin Diffusion Limit

 $C(\tau) \Box \exp(-\tau / \tau_c)$

 $J(\omega) \approx \frac{\tau_c}{1 + \omega^2 \tau_c^2}$

 $\omega \tau_c >> 1$

Optimal Control in Presence of Relaxation





$$\eta = \sqrt{1 + \xi^2} - \xi$$



The control problem



$$\frac{d}{dt} \begin{bmatrix} r_1 \\ r_2 \end{bmatrix} = \pi J \begin{bmatrix} -\xi u_1^2 & -u_1 u_2 \\ u_1 u_2 & -\xi u_2^2 \end{bmatrix} \begin{bmatrix} r_1 \\ r_2 \end{bmatrix}$$

Relaxation Optimized Pulse Elements (ROPE)



$$V(r_1, r_2) = \sqrt{\eta^2 r_1^2 + r_2^2} \qquad \eta = \sqrt{1 + \xi^2} - \xi$$

$$\frac{u_2 r_2}{u_1 r_1} = \eta$$

Comparison

Gain (ROPE/INEPT)



Relaxationoptimized pulse elements



Khaneja, Reiss, Luy, Glaser, J. Magn. Reson.162, 311 (2003)



Khaneja, Reiss, Luy, Glaser JMR(2003)

$$\frac{d}{dt} \begin{bmatrix} \langle I_z \rangle \\ \langle I_x \rangle \\ \langle 2I_y S_z \rangle \\ \langle 2I_z S_z \rangle \end{bmatrix} = \begin{bmatrix} 0 & -u(t) & & \\ u(t) & -k & -J & \\ & J & -k & -v(t) \\ & & v(t) & 0 \end{bmatrix} \begin{bmatrix} \langle I_z \rangle \\ \langle I_x \rangle \\ \langle 2I_y S_z \rangle \\ \langle 2I_z S_z \rangle \end{bmatrix}$$

Finite Horizon Problem



Khaneja, Reiss, Luy, Glaser JMR(2003)

Experimental Results



Khaneja, Reiss, Luy, Glaser JMR(2003)

Cross-Correlated Relaxation





Optimal control of spin dynamics in the presence of Cross-correlated Relaxation





Optimal trajectory preserves ratio $\frac{l_2}{l_1} = \eta$ and angle γ



$$\frac{l_2}{l_1} = \eta = \sqrt{1 + \xi^2} - \xi \qquad \qquad \xi = \sqrt{\frac{k_a^2 - k_c^2}{k_c^2 + J^2}}$$

Khaneja, Luy, Glaser PNAS(2003)



Experimental Transfer Functions



Khaneja, Luy, Glaser PNAS(2003)

TROPIC: Transverse relaxation optimized polarization transfer induced by cross-correlated relaxation.



Groel Protein: 800KDa

Room Temperature J = 93 HzProton Freq = 750Mhz Ka = 446 Hz

Kc = 326 Hz



Frueh et. al Journal of Biomolecular NMR (2005)



Control of Bloch Equations





Broadband Excitation







Inhomogeneous Ensemble of Bloch Equations

$$\frac{d}{dt}\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \begin{pmatrix} 0 & -\Delta \omega & -\varepsilon u(t) \\ \Delta \omega & 0 & -\varepsilon v(t) \\ \varepsilon u(t) & \varepsilon v(t) & 0 \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix}$$

$$g(\omega)$$

$$\omega \in [\omega_0 - B, \omega_0 + B]$$

$$\omega \in [\omega_0 - B, \omega_0 + B]$$



Robust Control Design



Robust Control Design by Area Generation

$$\frac{dX}{dt} = \varepsilon \left[u(t)\Omega_x + v(t)\Omega_y \right] X$$

$$U_{\varepsilon}(\Delta t) = \exp(-\varepsilon\Omega_y\Delta t) \exp(-\varepsilon\Omega_x\Delta t) \exp(\varepsilon\Omega_y\Delta t) \exp(\varepsilon\Omega_x\Delta t)$$

$$\approx I + (\Delta t)^2 \left[\underbrace{\varepsilon\Omega_x, \varepsilon\Omega_y}_{\varepsilon^2\Omega_z} \right]$$

$$U_{\varepsilon}(-\sqrt{\Delta t}) \exp(-\varepsilon\Omega_x\Delta t)U_{\varepsilon}(\sqrt{\Delta t}) \exp(\varepsilon\Omega_x\Delta t)$$

$$\approx I + (\Delta t)^2 \left[\underbrace{\varepsilon\Omega_x, \varepsilon\Omega_y}_{-\varepsilon^2\Omega_y} \right]$$

Lie Algebras, Areas and Robust Control Design

Using

f (ε)

3-8

<u>3-δ</u>

ng $\mathcal{E}\Omega_{y}, \mathcal{E}^{3}\Omega_{y}, \cdots, \mathcal{E}^{2k+1}\Omega_{y}$ as generators $f(\mathcal{E}) = \sum_{k} c_{k} \mathcal{E}^{2k+1}$ Choose $f(\mathcal{E})$ such that it is approx. constant for $\mathcal{E} \in [1 - \delta, 1 + \delta]$ $\exp(f(\mathcal{E})\Omega_{y})$

ε

 $\Theta(\varepsilon) = \exp(f_1(\varepsilon)\Omega_x)\exp(f_2(\varepsilon)\Omega_y)\exp(f_3(\varepsilon)\Omega_x)$

Fourier Synthesis Methods for Robust Control Design

$$U_{1} = \exp(k\pi\epsilon\Omega_{x})\exp(\epsilon\frac{\beta_{k}}{2}\Omega_{y})\exp(-k\pi\epsilon\Omega_{x})$$
$$= \exp(\epsilon\frac{\beta_{k}}{2}(\cos(k\pi\epsilon)\Omega_{y} + \sin(k\pi\epsilon)\Omega_{z}))$$

$$U_{2} = \exp(-k\pi\epsilon\Omega_{x})\exp(\epsilon\frac{\beta_{k}}{2}\Omega_{y})\exp(k\pi\epsilon\Omega_{x})$$
$$= \exp(\epsilon\frac{\beta_{k}}{2}(\cos(k\pi\epsilon)\Omega_{y} - \sin(k\pi\epsilon)\Omega_{z}))$$



$$\sum_{k} \beta_{k} \cos(k\pi\varepsilon) = \frac{\theta}{\varepsilon}$$





B. Pryor, N. Khaneja Journal of Chemical Physics (2006).

Fourier Synthesis Methods for Compensation



Time Optimal Control of Quantum Systems



Khaneja, Brockett, Glaser, Physical Review A, 63, 032308, 2001

Example

$$\dot{U} = -i \left(\begin{bmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_n \end{bmatrix} + \sum_i u_i B_i \right) U$$

 $B_i \in so(n), U \in \mathrm{SU}(\mathsf{n})$

Control Systems on Coset Spaces



Physical Review A , 63, 032308 (2001)

Cartan Decompositions, Two-Spin Systems and Canonical Decomposition of SU(4)



$$\frac{dU}{dt} = -i\left[\sum_{\alpha,\beta} J_{\alpha\beta} I_{\alpha} S_{\beta} + u_1 I_x + u_2 I_y + u_3 S_x + u_4 S_y\right] U$$

$$k = \left\{-i \ I_{\alpha}, -i S_{\beta}\right\}; \quad p = \left\{-i I_{\alpha} S_{\beta}\right\}$$

$$G = SU(4); \quad K = SU(2) \otimes SU(2)$$

$$I_{\alpha} = \sigma_{\alpha} \otimes I;$$

$$S_{\alpha} = I \otimes \sigma_{\alpha};$$

$$I_{\alpha}S_{\beta} = \sigma_{\alpha} \otimes \sigma_{\beta};$$

$$a = \{-iI_x S_x, -iI_y S_y, -iI_z S_z\}$$

$$G = K \exp(-i(\alpha_x I_x S_x + \alpha_y I_y S_y + \alpha_z I_z S_z)) K$$

$$\sigma_x = \frac{1}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}; \ \sigma_y = \frac{1}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}; \ \sigma_z = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Physical Review A , 63, 032308 (2001)

Geometry, Control and NMR

Indirect SWAP Operation

Efficiency η of indirect SWAP sequences



60 t [msec]

Reiss, Khaneja, Glaser J. Mag. Reson. 165 (2003)

Khaneja, et. al PRA(2007)



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